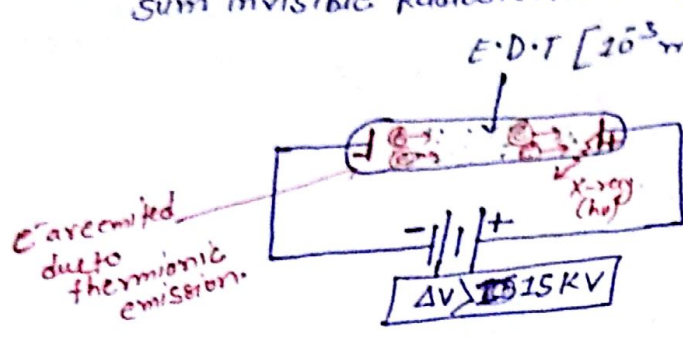


# 'X-RAYS'

Like photon.  $\rightarrow$  ~~is~~ undeviated from electric & magnetic field.

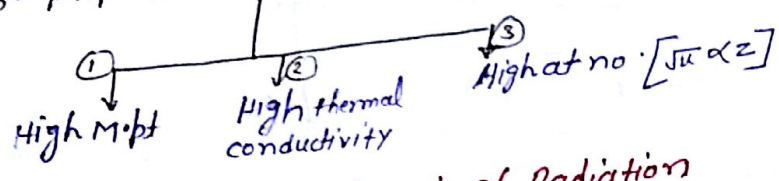
\* Reverse of P.E.E.  
 condn  $\rightarrow$  \* When highly energetic e's are made to strike metal target, electromagnetic radiation of the order  $0.01 \text{ \AA}$  to  $100 \text{ \AA}$  was observed & known as X-ray.

||  $\rightarrow$  Roentgen Exp:  $\rightarrow$  In  $\oplus$ nce of high potential & low potential, sum invisible radiation is emitted from anode which is called X-ray.

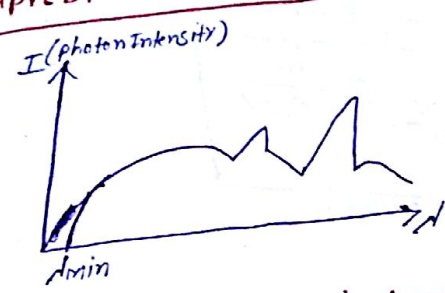


## Basic requirement of X-ray production

- 1a)  $\rightarrow$  e<sup>-</sup> producing source
- 1b)  $\rightarrow$  e<sup>-</sup> Accelerating source (more than 15KV pot.)
- 1c)  $\rightarrow$  Sp. property of target.



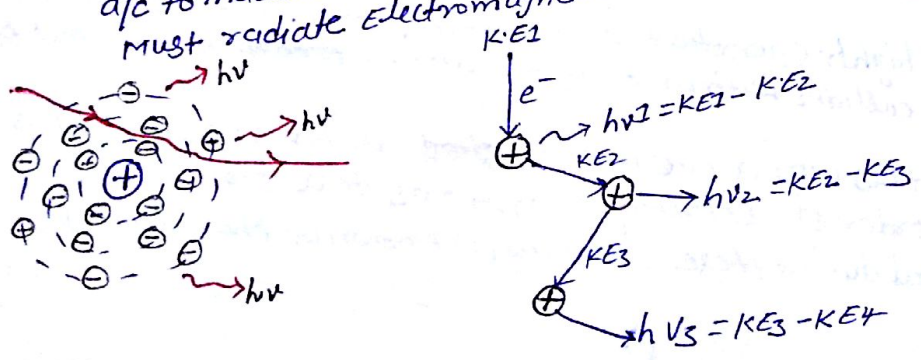
## # Graph b/w Intensity & Wavelength of Radiation



\* There was a wavelength, below which no radiation was obtained, this  $\lambda_{min}$  was named as cutoff wavelength.  
 \* If we observe the graph, then we see there are two type of variation, one is conti & the other is disconti or, in form of peak.

## # Production (Explanation) of continuous X-ray / White X-ray.

When highly energetic e's passed through an atom with very high K.E, they are strongly decelerated in electric field of nucleus & a/c to Maxwell electromagnetic theory every decelerating charge particle must radiate electromagnetic radiation & hence photon are emitted.



In successive collision  $e^-$  loses some part of its K.E energy & new photon with lost energy is generated.

\*\* If accelerating voltage across coolidge tube is ' $V_0$ '  
K.E of  $e^-$  just before hitting the target

$$\frac{1}{2} m v_e^2 = e V_0$$

$$v_e = \sqrt{\frac{2 e V_0}{m}}$$

\*\* Photon with highest energy is liberated when an  $e^-$  lose all its K.E in a single collision & corresponding to this loss only one photon come out.

$$h \nu_{\max} = e V_0 = K.E$$

$$\frac{h c}{\lambda_{\min}} = e V_0$$

$$\text{cutoff wave length} = \lambda_{\min} = \frac{h c}{e V_0}$$

\*\* other photons will have lesser energy so their wavelength will vary from  $\lambda_{\min}$  to  $\infty$ .

$$\Delta \cdot K_{\text{loss}} = E_{\text{x-ray}} = \frac{1}{2} m (v_1^2 - v_2^2)$$

$$\lambda_{\text{x-ray}} = \frac{h c}{E_{\text{x-ray}}} = \frac{h c}{\frac{1}{2} m (v_1^2 - v_2^2)}$$

\*\* Range of  $v_2 = 0 - \infty$

$$* v_2 - v_1 = E_{\text{x-ray}} = 0 \Rightarrow \lambda_{\max} = \infty$$

$$* v_2 = 0 \Rightarrow (E_{\text{x-ray}})_{\max} = \frac{1}{2} m v_1^2 = W = e \Delta V$$

$$\lambda_{\min} = \frac{h c}{(E_{\text{x-ray}})_{\max}} = \frac{h c}{\frac{1}{2} m v_1^2} = \frac{h c}{e \Delta V}$$

$$\lambda_{\min} = \frac{12400}{\Delta V (\text{volt})} \text{ \AA}$$

\*\*

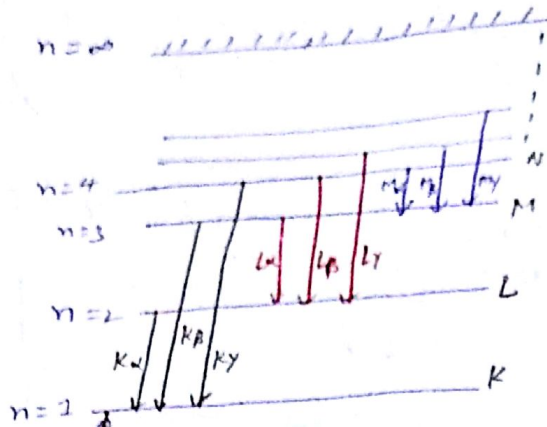
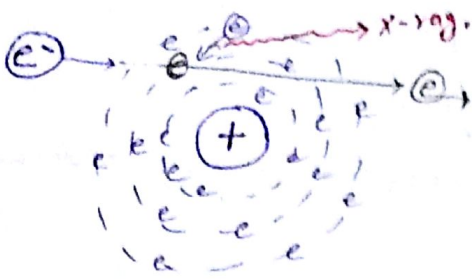
**NOTE** → Min wavelength of continuous X-ray only depend on anode potential. It is independent from at. no and nature of target.

2016  
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# Explanation or, production of characteristic X-ray or, explanation of peaks →

When highly energetic  $e^-$  enters into an atom, there is also some possibility that it collides with an  $e^-$  of the atom & knock out of the shell.

In this way a vacancy is ~~filled~~ created in shell, which is filled by higher order  $e^-$ s by making transitions to a lower stage, radiation emitted during these transition is known as characteristic X-ray.



$$V_{K\alpha} < V_{K\beta} < V_{K\gamma}$$

$$\lambda_{K\alpha} > \lambda_{K\beta} > \lambda_{K\gamma}$$

$$V_{L\alpha} > V_{L\beta} > V_{L\gamma}$$

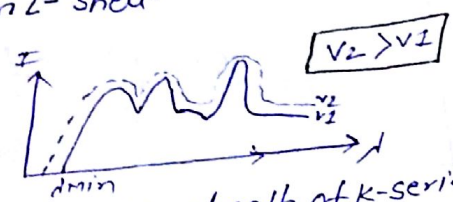
$$\lambda_{L\alpha} < \lambda_{L\beta} < \lambda_{L\gamma}$$

**NOTE** → Energy ( $E_K, E_L, E_M, E_N$ ) are not found from Bohr Model. If  $e^-$  is knocked out from L-shell.

$L\alpha, L\beta, L\gamma$

$$h\nu_{L\alpha} = E_L - E_M$$

$$h\nu_{L\beta} = E_L - E_N$$



- \* In a same transition ( $\alpha, \beta, \gamma$ ) Wavelength of  $\alpha$ -series is min.
- \* In a same series (K, L, M) Wavelength of  $\alpha$ -transition is max.
- \* Wavelength of characteristic X-ray doesn't depend on Accelerating voltage, they depend only on target material.

**Moseley's Law** → Square root of frequency  $\propto$  to effective atomic no. of multi  $e^-$  system.

bor 6  
↑  
Screening const.

$$\sqrt{\nu} \propto Z_{eff}$$

$$\sqrt{\nu} = aZ_{eff}$$

$$Z_{eff} = Z - b$$

At no.

- $K \Rightarrow b \Rightarrow 1$
- $L \Rightarrow b \Rightarrow 7.2$
- $M \Rightarrow b \Rightarrow 18.2$

**Screening const ( $\sigma$ )** → He explained screening const. as when an  $e^-$  whose a vacancy is created, then the effective nuclear charge  $\downarrow$  by a factor ' $\sigma$ ' known as screening const. In 'K' shell there is only one  $e^-$ , which belongs to 's'-orbitals i.e. spherically symmetry so for K-line, effective nuclear charge ~~become~~ become  $Z - 1$ .

\*\*

- \* For  $K\alpha$ -line:  $\sqrt{\nu_{K\alpha}} = \sqrt{\frac{3RC}{4}} (Z-1)$
- \* For  $K\beta$ -line:  $\sqrt{\nu_{K\beta}} = \sqrt{\frac{8RC}{9}} (Z-1)$
- \* For  $K\gamma$ -line:  $\sqrt{\nu_{K\gamma}} = \sqrt{\frac{15RC}{16}} (Z-1)$

\*\* → By this observation he concluded that property of material/atom depend on atomic no. not on atomic mass.

$$a = \text{proportional} = \sqrt{Rc \left( \frac{1}{n_1^2} - \frac{1}{n_2^2} \right)}$$

const

NOTE → 'a' depend on transition & 'b' depend on series but both are independent from at. no. of element.

$$\sqrt{V} = a(z-b)$$

$$V = R(z-b)^2 \left( \frac{1}{n_1^2} - \frac{1}{n_2^2} \right)$$

$$\frac{1}{\lambda} = R(z-b)^2 \left( \frac{1}{n_1^2} - \frac{1}{n_2^2} \right)$$

$$E = h\nu = Rch = (z-b)^2 \left( \frac{1}{n_1^2} - \frac{1}{n_2^2} \right)$$

\* R = Rydberg const =  $10^7 \text{ m}^{-1}$

\*  $\frac{1}{R} = 912 \text{ \AA}$

\* 1 Rydberg =  $Rch = 13.6 \text{ eV}$   
↑  
unit of energy

$$E = 13.6 (z-b)^2 \left( \frac{1}{n_1^2} - \frac{1}{n_2^2} \right) \text{ eV}$$

$$V = 2 \times 10^{25} (z-b)^2 \left( \frac{1}{n_1^2} - \frac{1}{n_2^2} \right)$$

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A → crystalline solid can cause x-ray to diffract.  
 R → Interatomic distance in crystalline solid is of order of 0.2 nm.

Ans → A